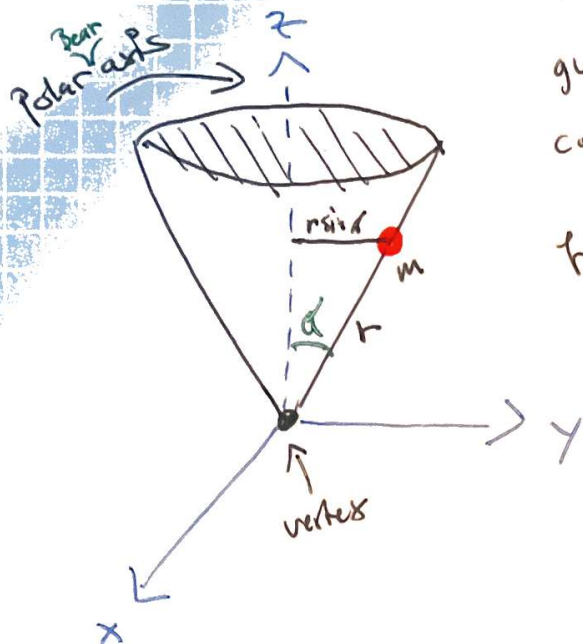


§14. Problem 2



I'm bored out of my mind, I guess this is ripe for spherical coordinates again!

$$\text{let } x = r \sin \alpha \cos \varphi$$

$$y = r \sin \alpha \sin \varphi$$

$$z = r \cos \alpha.$$

Note that α is fixed, so $\dot{\alpha} = 0$, hence the kinetic energy is $T = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\varphi}^2 \sin^2 \alpha)$. Taking the polar axis positive upward, $\vec{F} = -mg \hat{z}$. The force is connected to the potential energy by the relation

$$\vec{F} = - \frac{\partial U}{\partial \vec{r}} \text{ s.t. } U = - \int \vec{F} \cdot d\vec{r}. \text{ The upper limit}$$

in the integral can be chosen arbitrarily. If the force tends to zero at ∞ faster than $\frac{1}{r}$ then $\int_r^\infty F dr$ is convergent. Thus we put:

$$U(r) = \int_r^\infty F dr, \quad U(\infty) = 0. \text{ The choice of}$$

arbitrary constant here is called the gauge. Anywho,

$$U(r) = - \int_z^\infty mg dz = mgz.$$

The Lagrangian is:

$$\mathcal{L} = \frac{1}{2} m (\dot{r}^2 + r^2 \dot{\varphi}^2 \sin^2 \alpha) - mgr \cos \alpha$$

Noting through repetition, the cyclic coordinate is φ , hence $p_\varphi = \text{const}$. This means that

$$p_\varphi = \frac{\partial \mathcal{L}}{\partial \dot{\varphi}} = m r^2 \dot{\varphi} \sin^2 \alpha \Leftrightarrow \dot{\varphi} = \frac{p_\varphi}{m r^2 \sin^2 \alpha}$$

$\mathcal{L}$ has no explicit t so energy is conserved.

$$\begin{aligned} E &= \frac{1}{2} m \dot{r}^2 + \frac{1}{2} m r^2 \dot{\varphi}^2 \sin^2 \alpha + mgr \cos \alpha \\ &= \frac{1}{2} m \dot{r}^2 + \underbrace{\frac{p_\varphi^2}{2 m r^2 \sin^2 \alpha}}_{= U_{\text{eff}}(r)} + mgr \cos \alpha \end{aligned}$$

Back to 1D business, we can solve for r by:

$$t = \int \frac{dr}{\sqrt{\frac{2}{m} [E - U_{\text{eff}}(r)]}} \quad \text{and}$$

$$\varphi = \frac{p_\varphi}{\sqrt{2m} \sin^2 \alpha} \int \frac{dr}{r^2 \sqrt{E - U_{\text{eff}}(r)}}$$

Interchange $M_\varphi = p_\varphi$ as you like, nonetheless, the standard answer.